BULLETIN OF THE CHEMICAL SOCIETY OF JAPAN, VOL. 51 (6), 1885—1886 (1978)

## Location of the $\nu_1$ , $\nu_2$ , $\nu_3$ , and $\nu_5$ Fundamental Vibrations of Ethylene- $d_4$

Nobukimi Ohashi, Tetsuo Kyogoku, and Shin-ichi Matsuoka

Institute for Spectroscopic Study of Matter, Faculty of Science, Kanazawa University, Maruno-uchi, Kanazawa 920 (Received November 15, 1977)

**Synopsis.** The  $^{Q}Q$  branch peaks of the  $v_1$ ,  $v_2$ , and  $v_3$  Raman active bands of ethylene- $d_4$  were located at  $v_1$ =  $2260.8\pm0.1~\mathrm{cm^{-1}}$ ,  $v_2$ = $1518.3\pm0.1~\mathrm{cm^{-1}}$ , and  $v_3$ = $984.4\pm0.1~\mathrm{cm^{-1}}$ , respectively. The band center of the  $v_5$  band which does not have a  $^{Q}Q$  branch was determined to be  $v_5$ = $2315.6\pm0.3~\mathrm{cm^{-1}}$  from the analysis in the symmetric top approximation.

While many Raman data have been presented for ethylene- $d_0$ ,  $^{1-6}$ ) few Raman data have been published for ethylene- $d_4$ . Though its pure rotational Raman spectrum was studied by Dowling et al.,  $^2$ ) its vibrotational Raman spectra have not been studied. Only the liquid state vibrational frequencies of the Raman active bands have been measured, by Hemptinne et al.  $^7$ ) As a preliminary to the vib-rotational Raman study of ethylene- $d_4$ , we report the  $v_1$ ,  $v_2$ ,  $v_3$ , and  $v_5$  fundamental frequencies, since these data are considered to be important in the force constant calculation of this molecule.

## **Experimental**

The experimental setup consisted of a He-Cd laser operated at 441.570 nm, a gas cell, and an Echelle grating monochromator. The scattered Raman radiation from a multi-reflection cell<sup>8)</sup> placed inside the laser cavity was analyzed by the Echelle grating monochromator and detected by a cooled photomultiplier operated in conjunction with photon counting equipment. Ethylene- $d_4$  gas obtained from Merck, Sharp, and Dohme was used at a pressure of about 400 Torr. Wavelength calibration was made by using as standards the lines from an iron-neon hollow cathode lamp. A slit width of 0.4 cm<sup>-1</sup> was used in the measurements of the  $v_1$   ${}^{Q}Q$  and the  $v_2$   ${}^{Q}Q$  branch peaks, and a

slit width of  $0.8 \text{ cm}^{-1}$  was used in the measurements of the  $v_3$  Q branch peak and the  $v_5$  band.

## Results

The  $^{\rm Q}{\rm Q}$  branch peaks of the totally symmetric bands of ethylene- $d_4$  gas were determined to be  $v_1=2260.8\pm0.1$  cm<sup>-1</sup>,  $v_2=1518.3\pm0.1$  cm<sup>-1</sup>, and  $v_3=984.4\pm0.1$  cm<sup>-1</sup>. The  $v_1$  value is larger than that in liquid state by 10 cm<sup>-1</sup>, but the  $v_2$  and  $v_3$  fundamental frequencies are nearly equal to those in liquid state, as obtained by Hemptinne et al.<sup>7</sup> These circumstances are similar to those in ethylene- $d_0$ .<sup>9</sup>

As the rotational structure of the  $v_5$  band was observed under low resolution (about 1.0 cm<sup>-1</sup>), as shown in Fig. 1, this band was analyzed in the symmetric top approximation. In order to enable the K-assignment to be done exactly we made use of the simulated spectrum which was caculated using the calculation program for the asymmetric rotor. As the rotational constants in the  $v_5=1$  state can not be obtained at the present stage, we calculated the spectrum with  $A_0=A_5=2.425$ cm<sup>-1</sup>,  $B_0 = B_5 = 0.73630 \text{ cm}^{-1}$ , and  $C_0 = C_5 = 0.56355$ cm $^{-1}$ , where the ground state constants for ethylene- $d_4$ were taken from the paper of Duncan et al. 10) In the symmetric top approximation the P,RQk branches form a series of lines at the origins of subbands, if  $\bar{B}_0 - \bar{B}_5$  is negligibly small. From the least squares fit of the observed data we obtained the following quadratic equation for the subband origins:

$$v^{\text{sub}} = 2317.4 + 3.601K - 0.006K^2$$
.

This equation gives  $v_5 = 2315.6 \pm 0.3$  cm<sup>-1</sup> for the band

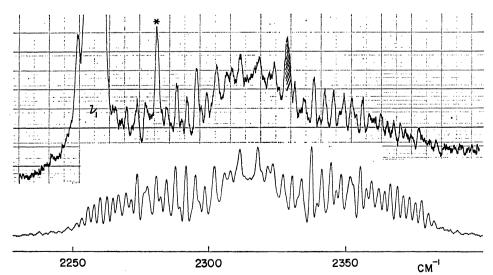


Fig. 1. Observed (top) and calculated (bottom) spectra of the  $\nu_5$  band of ethylene- $d_4$ . The peak marked by \* was attributed to  $\nu_1$  (C<sub>2</sub>HD<sub>3</sub>) and the peak with shadow was overlapped by N<sub>2</sub> line.

center, which is also about  $11~\rm cm^{-1}$  larger than that in liquid state. The negative value of the coefficient in  $K^2$  is contradictory to  $A_5 - \bar{B}_5 = 1.801 > A_0 - \bar{B}_0 = 1.775$ . It was perhaps caused by the circumstance that the correction for the term  $(\bar{B}_5 - \bar{B}_0)J(J+1)$  was not applied to the observed data and that the symmetric top approximation was used for an asymmetric top molecule.

## References

- 1) J. Romanko, T. Feldman, E. J. Stransbury, and A. McKellar, Can. J. Phys., 32, 735 (1954).
- 2) J. M. Dowling and B. P. Stoicheff, Can. J. Phys., 37, 703 (1959).
  - 3) T. Feldman, J. Romanko, and H. L. Welsh, Can. J.

Phys., 34, 737 (1956).

- 4) L. Nemes and S. Suzuki, *J. Raman Spectrosc.*, 2, 193 (1974).
- 5) G. W. Hills and W. J. Jones, J. Chem. Soc., Faraday Trans. 2, 71, 812 (1975).
- 6) R. B. Foster, G. W. Hills, and W. J. Jones, *Mol. Phys.*, **33**, 1589 (1977).
- 7) M. de Hemptinne, J. Jungers, and J. M. Delfosse, J. Chem. Phys., 6, 319 (1938).
- 8) N. Ohashi, H. Watanabe, and S. Matsuoka, *Jpn. Appl. Phys.*, **12**, 1103 (1973).
- 9) G. Glocker and M. M. Renfrew, J. Chem. Phys., 6, 170 (1938); 6, 409 (1938).
- 10) J. L. Duncan, I. J. Wright, and D. Van Lerberghe, *J. Mol. Spectrosc.*, **42**, 463 (1972).